# Infinite Particle Systems in Non-Discrete Spaces: Orthogonal Intertwiners

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= a time-continuous Markov process describing the stochastic evolution of particles in a set E given a deterministic initial configuration

Uniform sticky Brownian motions with stickiness  $\theta > 0$ Special case of a martingale problem introduced by Howitt, Warren '09.

Let  $E = \mathbb{R}$ .



# Different notations

Unlabeled notation	Labeled notation
Markov process	family of Markov processes indexed via number of particles $N \in \mathbb{N}$
values in <b>N</b> , the set of counting measures consisting of $\mu = \sum_{k=1}^{N} \delta_{x_k}$ , $x_k \in \mathbb{R}$ . $\mu(A) =$ number of particles in a subset $A$	values in $\mathbb{R}^N$
$(P_t)_{t\geq 0}$ Markov semigroup acting on functions $F: \mathbf{N}  ightarrow \mathbb{R}$	family of Markov semigroups $(P_t^{[N]})_{t\geq 0}, N\in \mathbb{N}$ acting on functions $f_N: \mathbb{R}^N  o \mathbb{R}$

## n-th factorial measure intertwiner

For  $f_n : \mathbb{R}^n \to \mathbb{R}$  and a configuration  $\sum_{k=1}^N \delta_{x_k}$  put

$$J_n\left(f_n, \sum_{k=1}^N \delta_{x_k}\right) := \sum_{\substack{i_1, \dots, i_n = 1 \\ \text{pairwise different}}}^N f_n(x_{i_1}, \dots, x_{i_n})$$

$$J_1(f_1, \delta_x + \delta_y + \delta_z) = f_1(x) + f_1(y) + f_1(z),$$
  

$$J_2(f_2, \delta_x + \delta_y + \delta_z) = f_2(x, y) + f_2(y, x) + f_2(x, z) + f_2(z, x) + f_2(y, z) + f_2(z, y)$$
  
Theorem 1 (W., '23)

 $J_n$  intertwines the dynamics of infinitely many sticky Brownian motions and their *n*-particle evolution, *i.e.*,

$$P_{t}J_{n}(f_{n},\cdot)(\mu) = J_{n}\left(P_{t}^{[n]}f_{n},\mu\right) \qquad f_{n}:\mathbb{R}^{n}\to\mathbb{R}, \quad \mu\in\mathbb{N}$$

### Example

Given an initial configuration  $\sum_{k=1}^{N} \delta_{x_k}$ : What is the expected number of particles in a set A at a time t > 0?

expected number of particles in a set A at time t

$$= P_t (\nu \mapsto \nu(A)) \left( \sum_{k=1}^N \delta_{x_k} \right)$$
$$= P_t J_1(\mathbf{1}_A, \cdot) \left( \sum_{k=1}^N \delta_{x_k} \right)$$
$$= J_1 \left( P_t^{[1]} \mathbf{1}_A, \mu \right)$$
$$= \sum_{k=1}^N P_t^{[1]} \mathbf{1}_A(x_k) =$$

Let each particle evolve the one-particle dynamics separately. Sum of the probabilities that the respective particle is in the set A at time t.

The proof relies solely on compatibility (Le Jan, Raimond, '04)

= "the action of removing a particle commutes with the dynamics".

Le Jan, Raimond proved a one-to-one correspondence to stochastic flows.

#### **Examples:**

- independent particles:
  - independent random walkers (IRW)
  - free Kawasaki dynamics
- correlated Brownian motions
- coalescing Brownian motions

#### Theorem 2 (Redig, Jansen, Floreani, W., '21)

Consider a finite particle system that preserves the number of particles. Then,

consistency  $\iff J_n \text{ intertwines } P_t^{[n]} \text{ and } P_t \text{ for all } t \text{ and } n, \text{ i.e.,}$  $P_t J_n(f_n, \cdot)(\mu) = J_n\left(P_t^{[n]}f_n, \mu\right), \qquad f_n : \mathbb{R}^n \to \mathbb{R}, \quad \mu \in \mathbf{N}$ 

consistency = "the action of removing a particle **uniformly at random** commutes with the dynamics".

#### Examples:

- symmetric inclusion process (SIP), symmetric exclusion process (SEP).
  - Theorem 2 recovers well-known self-dualities in terms of falling factorial polynomials.
- generalized SIP in the continuum = Moran process from population genetics

# Infinite-dimensional orthogonal polynomials

Let  $\rho$  be a probability measure on  $\mathbf{N},$  i.e., the distribution of a point process.

orthogonal projection of 
$$J_n(f_n, \cdot)$$
  
 $I_n(f_n, \cdot) :=$ onto  $\{J_k(u_k, \cdot), u_k : \mathbb{R}^k \to \mathbb{R}, 0 \le k \le n-1\}^{\perp}$   
in  $L^2(\rho)$ .

Keywords: infinite-dimensional orthogonal polynomials, (extended) Fock spaces, chaos decompositions, multiple stochastic integrals, non-Gaussian white noise analysis, Malliavin calculus

Let  $\rho$  be the distribution of a Pascal process with parameters  $\theta Vol_{\mathbb{R}}$  and arbitrary  $\rho \in (0, 1)$ .

## Theorem 3 (W., '23)

 $I_n$  intertwines the dynamics of infinitely many sticky Brownian motions and their *n*-particle evolution, i.e.,

$$P_t I_n(f_n, \cdot)(\mu) = I_n\left(P_t^{[n]} f_n, \mu\right) \qquad f_n : \mathbb{R}^n \to \mathbb{R}, \quad \mu \in \mathbf{N}$$

## Corollary 4 (W., '23)

 $\rho$  is reversible for infinitely many sticky Brownian motions.

The proof of the corollary applies a result of Brockington, Warren '23: existence of a reversible measure for *n* ordered uniform sticky Brownian motions for each  $n \in \mathbb{N}$ .

Connection: reversible measures and orthogonal intertwiners Theorem 5 (Redig, Jansen, Floreani, W., '21)

$$J_n$$
 intertwines  $P_t^{[n]}$  and  $P_t$  for all t and  $n + \rho$  reversible  
 $\downarrow \downarrow$   
 $I_n$  intertwines  $P_t^{[n]}$  and  $P_t$  for all t and r

Together with Theorem 2:

finite number of particles +  $\rho$  reversible + consistency  $\Rightarrow$   $I_n$  intertwines  $P_t^{[n]}$  and  $P_t$ 

Applications:

- Theorem 5 recovers well-known self-dualities in terms of orthogonal polynomials for discrete particle systems (SIP, SEP, IRW)
- independent particles ( $\rho$  distribution of the Poisson process)
- generalized SIP ( $\rho$  distribution of the Pascal process)

# Thank you!

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