# Intertwining for continuum interacting particle systems 

Workshop: Population Genetics, Interacting Particle Systems and Stochastic Flows: a duality perspective

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## Overview

## Joint work with

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1. Known facts - discrete spaces.

Self-dualities of Markov processes describing the evolution of particles on a discrete set.
2. My Research - general spaces.

What happens if we replace the discrete space by a much more general space?

## Duality functions

## Definition (Stochastic duality of Markov processes)

Let $X=\left(\Omega_{1}, \mathcal{F}_{1},\left(X_{t}\right)_{t \geq 0},\left(\mathbb{P}_{x}\right)_{x \in \mathbb{X}}\right)$ and $Y=\left(\Omega_{2}, \mathcal{F}_{2},\left(Y_{t}\right)_{t \geq 0},\left(\mathbb{P}^{y}\right)_{y \in \mathbb{Y}}\right)$ be two (time-continuous) Markov processes with state spaces $\mathbb{X}, \mathbb{Y} . X$ and $Y$ are dual with respect to $H: \mathbb{X} \times \mathbb{Y} \rightarrow \mathbb{R}$ if and only if for all $x \in \mathbb{X}, y \in \mathbb{X}$ and $t \geq 0$

$$
\mathbb{E}_{x} H\left(X_{t}, y\right)=\mathbb{E}^{y} H\left(x, Y_{t}\right)
$$

Semigroup notation: $P_{t} H(\cdot, y)(x)=S_{t} H(x, \cdot)(y)$.

We consider self-duality, i.e, $X=Y$.

## Self-Duality for SIP (falling factorials)

Consider the symmetric Inclusion Process (SIP) generated by

$$
L f(x)=\sum_{k \in E} \sum_{\ell \in E}\left(f\left(x-\delta_{k}+\delta_{\ell}\right)-f(x)\right) c(k, \ell)\left(\alpha_{\ell}+x_{\ell}\right) x_{k}
$$

on the configuration space $\mathbb{N}_{0}^{E}$ with finite $E=\{1, \ldots, N\}$ and symmetric conductances $c(k, \ell)=c(\ell, k) \geq 0, k, \ell \in E$ and $\alpha=\left(\alpha_{1}, \ldots, \alpha_{N}\right), \alpha_{k} \geq 0$. A reversible measure (i.e. detailed balance) is given by

$$
\rho_{\alpha, p}=\bigotimes_{k \in E} \operatorname{NegativeBinomial}\left(p, \alpha_{k}\right)
$$

for each $p \in(0,1)$.

Similarly:

- independent random walkers (IRW): Product of Poisson distributions
- symmetric exclusion process (SEP): Product of Binomial distributions

Theorem (Carinci, Giardinà, Redig, '19)
Let $\rho$ be the reversible measure, $(n)_{k}:=n(n-1) \cdots(n-k+1)$. Then, a self-duality function for IRW, SIP, SEP is

$$
H(x, y):=\frac{1}{\rho(\{x\})} \prod_{k \in E} \frac{1}{x_{k}!}\left(y_{k}\right)_{x_{k}} .
$$

## Orthogonal dualities

Let $\left(M_{n}(\cdot, a, p)\right)_{n \in \mathbb{N}_{0}}$ be the (monic) Meixner polynomials (orthogonal with respect to NegativeBinomial $(p, a)$. Consider the multivariate orthogonal polynomials

$$
P_{x}(y, \alpha):=\frac{1}{\rho(\{x\})} \prod_{k \in E} \frac{1}{x_{k}!} P_{x_{k}}\left(y_{k}, \alpha_{k}\right)
$$

Theorem (Franceschini, Giardinà, 19')
$H(x, y):=P_{x}(y, \alpha)$ is a self-duality function for the SIP.
IRW, SEP similarly.

## Consistency

"the action of removing a particle uniformly at random commutes with the dynamic"

SIP, SEP, IRW share the following relation: $\mathcal{A} P_{t}=P_{t} \mathcal{A}$ for the so-called lowering operator $\mathcal{A} f(x):=\sum_{k \in E} x_{k} f\left(x-\delta_{k}\right)$.
In terms of expectations, for each $f: \mathbb{N}_{0}^{E} \rightarrow \mathbb{R}, t \geq 0, x \in \mathbb{N}_{0}^{E}$

$$
\mathbb{E}_{X}\left[\sum_{k \in E} f\left(X_{t}-\delta_{k}\right) X_{t}\right]=\sum_{k \in E} x_{k} \mathbb{E}_{x-\delta_{k}}\left[f\left(X_{t}\right)\right]
$$

For reversible particle systems on discrete sets, the property of consistency is equivalent to self-duality (Carinci, Giardinà, Redig, '19).

## Non-discrete spaces?

Question: How to generalize these dualities to the continuum? More precisely: Replace discrete $E$ by $\mathbb{R}$.
Challenges: How to generalize ...

1. the configuration spaces?
2. the models?
3. consistency?
4. the concept of duality functions?
5. falling factorials?
6. reversible measures?
7. orthogonal polynomials?

Further challenges:
8. Algebraic properties?
9. Infinitely many particles?

## Main Idea

SIP, SEP, IRW are consistent. The concept of consistency can be generalized naturally. It turns out that this is the right notation and starting point to obtain dualities.

## Configurations

Let $E=\mathbb{R}$. Model a configuration as counting measures, i.e,

$$
\mathbf{N}:=\left\{\sum_{k=1}^{n} \delta_{x_{k}}: x_{k} \in E, n \in \mathbb{N}_{0} \cup\{\infty\}\right\}
$$

Modern notation for Point processes (Last / Penrose).

## Models in the continuum

- generalized Version of the SIP on the continuum (gSIP): Let $\alpha$ be a finite measure on $\mathbb{R}$ and $c: \mathbb{R} \times \mathbb{R} \rightarrow[0, \infty)$ symmetric. Consider the process generated by

$$
L f(\eta)=\iint c(x, y)\left(f\left(\eta-\delta_{x}+\delta_{y}\right)-f(\eta)\right)(\eta+\alpha)(\mathrm{d} y) \eta(\mathrm{d} x)
$$

for $f: \mathbf{N} \rightarrow \mathbb{R}, \eta \in \mathbf{N}$.

- Independent Markov processes, e.g., free Kawasaki
- Strongly Consistent systems (Called compatibility by Le Jan, Raimond) stochastic flows
- Brownian motions
- Correlated Brownian motions


## Consistency

Define

$$
\mathcal{A} f(\eta):=\int f\left(\eta-\delta_{x}\right) \eta(\mathrm{d} x)
$$

We say that a Markov process is consistent if

$$
\mathcal{A} P_{t}=\mathcal{A} P_{t}
$$

i.e. for each $f: \mathbf{N} \rightarrow \mathbb{R}, t \geq 0, \eta \in \mathbf{N}$ In other words,

$$
\mathbb{E}_{\eta}\left[\int f\left(\eta_{t}-\delta_{x}\right) \eta_{t}(\mathrm{~d} x)\right]=\int \mathbb{E}_{\eta-\delta_{x}}\left[f\left(\eta_{t}\right)\right] \eta(\mathrm{d} x)
$$

## Duality and Intertwiners

Let $E$ be discrete. Let $\rho$ be a reversible measure for $X$ and $H$ is a self-duality function. Put a linear operator

$$
T f(y)=\int H(x, y) f(x) \rho(\mathrm{d} x), \quad y \in E
$$

for functions $f: E \rightarrow \mathbb{R}$. Then, $T$ intertwines $P_{t}$ with itself.

## Example: intertwiners for discrete systems

- Duality $\frac{1}{\rho(\{x\})} \prod_{k \in E} \frac{1}{x_{k}!}\left(y_{k}\right)_{x_{k}}$ with falling factorials turns into

$$
T f(y)=\int H(x, y) f(x) \rho(\mathrm{d} x)=\sum_{x_{1}=0}^{y_{1}} \cdots \sum_{x_{N}=0}^{y_{N}}\binom{y_{1}}{x_{1}} \cdots\binom{y_{N}}{x_{N}} f(x)
$$

- Duality $\frac{1}{\rho(\{x\})} \prod_{k \in E} \frac{1}{x_{k}!} P_{x_{k}}\left(y_{k}, \alpha_{k}\right)$ with orthogonal polynomials turns into

$$
T f(y)=\int H(x, y) f(x) \rho(\mathrm{d} x)=\sum_{x_{1}=0}^{\infty} \cdots \sum_{x_{N}=0}^{\infty} \frac{1}{x_{1}!} P_{x_{1}}\left(y_{1}, \alpha_{1}\right) \cdots \frac{1}{x_{N}!} P_{x_{N}}\left(y_{N}, \alpha_{N}\right) f(x)
$$

## Improve duality functions

Moreover, if $S$ intertwines $P_{t}$ with itself, and $H$ is a self-duality function (for example the cheap-duality function $\left.H(x, y)=\delta_{x, y} \rho(\{x\})\right)$, then $\tilde{H}(x, y)=T H(\cdot, y)(x)$ is another self-duality function.

## Why intertwiners?

- Intertwiners can be seen as a "lifting" of duality functions.
- The intertwiners can be generalized naturally
- These generalization lead to kernel operators without absolute continuity with respect to a reversible measure.


## Generalized falling factorials

The generalization does already exist: Lennard's K-transform - special kind of Möbius transform. Put

$$
K f(\eta)=\sum_{\nu \hat{\leq} \eta, \nu(\mathbb{R})<\infty} f(\nu)
$$

$\hat{\leq}$ means that e.g. $\nu=\delta_{x}$ occurs two times if $\eta=2 \delta_{x}$. Other representation:

$$
K f(\eta)=\sum_{n=0}^{\infty} \frac{1}{n!} \int f\left(\delta_{x_{1}}+\ldots+\delta_{x_{n}}\right) \eta^{(n)}\left(\mathrm{d}\left(x_{1}, \ldots, x_{n}\right)\right)
$$

$\eta^{(n)}$ falling factorial measure, can be seen as a generalization of falling factorials.

## Link to the discrete world

Theorem (Redig, Jansen, Floreani, W., '21)
Let $X$ be consistent. Then $K P_{t}=P_{t} K$.

Moreover, if $X$ conserves the number of particles then
$\mathbb{E}_{\eta}\left[\int f\left(\delta_{y_{1}}+\ldots+\delta_{y_{n}}\right) \eta_{t}^{(n)}\left(\mathrm{d}\left(y_{1}, \ldots, y_{n}\right)\right)\right]=\int \mathbb{E}_{\delta_{x_{1}}+\ldots+\delta_{x_{n}}}\left[f\left(\eta_{t}\right)\right] \eta^{(n)}\left(\mathrm{d}\left(x_{1}, \ldots, x_{n}\right)\right)$
Intertwines the dynamics of arbitrary many particles with the dynamic of only $n$-particles.

Intertwining for all consistent systems (not only SIP).

## Recover self-duality of discrete systems

If $f(\eta)=\mathbb{1}_{\eta\left(D_{1}\right)=d_{1}} \cdots \mathbb{1}_{\eta\left(D_{N}\right)=d_{N}}, D_{1}, \ldots, D_{N}$ partition of $\mathbb{R}, d_{1}, \ldots, d_{N} \in \mathbb{N}_{0}$, then,

$$
K f(\eta)=\prod_{k=1}^{N} \frac{\left(\eta\left(D_{k}\right)\right)_{d_{k}}}{d_{k}!}
$$

In particular, if $E$ is finite, then $K=T$.

## How to generalize reversible measures?

Observation: Let $E$ be finite, $X \sim \rho_{\alpha, p}$ and write

$$
X_{D}:=\sum_{k \in D} x_{k}, \quad D \subset E .
$$

Then:

1. $X_{D} \sim$ NegativeBinomial $\left(\sum_{k \in D} \alpha_{k}, p\right)$ (Negative Binomial distribution forms a convolution semigroup);
2. $X_{D}$ and $X_{D \prime}$ are independent for disjoint $D, D^{\prime}$.

We look for a measure $\rho$ on $\mathbf{N}$. If $\rho$ is a probability measure, $\rho$ is the distribution of a point process.

Question: Is for a measure $\alpha$ on $\mathbb{R}$ a point process $\zeta$ with

1. $\xi\left(A_{1}\right), \ldots, \xi\left(A_{N}\right)$ are independent for pairwise disjoint measurable sets $A_{1}, \ldots, A_{n} \subset \mathbb{R}$,
2. $\xi(A) \sim \operatorname{NegativeBinomial}(p, \alpha(A))$ for each measurable $A \subset \mathbb{R}$

Yes: theory of Lévy processes on general spaces (in our examples: Pascal process).

Indeed: The distribution of the Pascal process is reversible for the gSIP.

## Infinite dimensional polynomials

Define polynomials $\mathcal{P}_{n}$ of degree $\leq n$ as the linear combinations of the functions

$$
\eta \mapsto \int f_{k} \mathrm{~d} \eta^{\otimes k}
$$

with $k \leq n, f_{k}: \mathbb{R}^{k} \rightarrow \mathbb{R}$.

Motivation if is $E$ is discrete. $\eta=x_{1} \delta_{1}+\ldots+x_{N} \delta_{N}$, then

$$
\int f_{k} \mathrm{~d} \eta^{\otimes k}=\sum_{\ell_{1}, \ldots, \ell_{n}=1}^{N} f_{k}\left(\ell_{1}, \ldots, \ell_{n}\right) x_{\ell_{1}} \cdots x_{\ell_{n}}
$$

is a multivariate polynomial of degree $n . x_{1}^{n}=\int f_{k} \mathrm{~d} \eta^{\otimes n}$ can be recovered by

$$
f_{k}\left(\ell_{1}, \ldots, \ell_{n}\right)= \begin{cases}1 & \ell_{1}=\cdots=\ell_{n}=1 \\ 0 & \text { else }\end{cases}
$$

## Orthogonal polynomials

For $f_{n}: \mathbb{R}^{n} \rightarrow \mathbb{R}$ put

$$
I_{n} f_{n}=\text { orthogonal projection of }\left(\eta \mapsto \int f_{n} \mathrm{~d} \eta^{\otimes n}\right) \text { onto } \mathcal{P}_{n-1}^{\perp} \text { in } L^{2}(\rho)
$$

$\rho$ reversible measure (e.g. of the Pascal process).

Studied by e.g. by Lytvynov, '02—Link to infinite dimensional analysis, chaos decompositions and multiple stochastic integrals. Poisson case: multiple Wiener Itô integrals.

## Factorization property

## Theorem (Redig, Jansen, Floreani, W., '21)

Suppose that $\rho$ is the distribution of some finite completely independent point process. Let $N \geq 2, A_{1}, \ldots, A_{N} \subset \mathbb{R}$, measurable, pairwise disjoint, and $d_{1}, \ldots, d_{N} \in \mathbb{N}_{0}$. Further let $f_{i}: \mathbb{R}^{d_{i}} \rightarrow \mathbb{R}, i=1, \ldots, N$ be bounded measurable functions that vanish on $\mathbb{R}^{d_{i}} \backslash A_{i}^{d_{i}}$. Set $n:=d_{1}+\cdots+d_{N}$. Then

$$
I_{n}\left(f_{1} \otimes \ldots \otimes f_{n}\right)(\eta)=I_{d_{1}} f_{1}(\eta) \cdots I_{d_{n}} f_{n}(\eta)
$$

for $\rho$-almost all $\eta \in \mathbf{N}_{<\infty}$.

## Properties of infinite dimensional Meixner Polynomials

Theorem (Redig, Jansen, Floreani, W., '21)

- $I_{d} \mathbb{1}_{A^{d}}(\eta)=M_{d}(\eta(A) ; \alpha(A) ; p)$
- There are measures $\lambda_{n}$ such that

$$
\int\left(I_{n} f_{n}\right)\left(I_{m} g_{m}\right) \mathrm{d} \rho=\mathbb{1}_{\{n=m\}} \int f_{n} g_{m} \mathrm{~d} \lambda_{n}
$$

for permutation invariant $f_{n}: \mathbb{R}^{n} \rightarrow \mathbb{R}, g_{m}: \mathbb{R}^{m} \rightarrow \mathbb{R}, n, m \in \mathbb{N}_{0}$.

## Intertwining

## Theorem (Redig, Jansen, Floreani, W., '21)

Assume that $X$ is consistent and conserves the number of particles (then the n-particle semigroup $P_{t}^{[n]}$ acting on permutation invariant functions $f_{n}: E^{n} \rightarrow \mathbb{R}$ is well-defined). Assume that $\rho$ is reversible. Then

$$
P_{t} I_{n}=I_{n} P_{t}^{[n]}
$$

For $f: \mathbf{N} \rightarrow \mathbb{R}$ put

$$
T f=\sum_{n=0}^{\infty} \frac{1}{n!} I_{n} f_{n}
$$

by $f\left(\delta_{x_{1}}+\ldots+\delta_{x_{n}}\right)=f_{n}\left(x_{1}, \ldots, x_{n}\right)$. Then, $P_{t} T=T P_{t}$.

## Recover discrete dualities

Theorem (Redig, Jansen, Floreani, W., '21)
$T$ recovers duality functions for SIP (also IRW)!

Indeed, if $f(\eta)=\mathbb{1}_{\eta\left(D_{1}\right)=d_{1}} \cdots \mathbb{1}_{\eta\left(D_{N}\right)=d_{N}}, D_{1}, \ldots, D_{N}$ partition of $\mathbb{R}, d_{1}, \ldots, d_{N} \in \mathbb{N}_{0}$, then,

$$
T f(\eta)=\prod_{k=1}^{N} \frac{M_{d_{k}}\left(\eta\left(D_{k}\right) ; \alpha\left(D_{k}\right) ; p\right)}{d_{k}!}
$$

## Further challenges

- Duality and Lie-Algebra representations
- Infinitely many particles


## Thank you!

S. Floreani, S. Jansen, F. Redig, S.W.: Duality and intertwining for consistent Markov processes arXiv:2112.11885 [math.PR], 32 pp.

