

INTRODUCTION TO BOHMIAN MECHANICS
SUMMER TERM 2016

EXERCISE SHEET 2

Exercise 1: Lagrangian Formalism

This exercise deals with the Lagrangian formalism of classical mechanics.

The Lagrange function is a function of coordinates and velocities (thought of as velocities of trajectories passing through those coordinates): $\mathcal{L} = \mathcal{L}(q, \dot{q})$, where $q = (q_1, \dots, q_N)$ is an element of the configuration space \mathbb{R}^N .

- a) Integrating the Lagrangian along any curve $\gamma : [t_0, t] \rightarrow \mathbb{R}^N$ gives the action of that path:

$$S[\gamma] = \int_{t_0}^t \mathcal{L}(\gamma(s), \dot{\gamma}(s)) ds.$$

Show, without full rigor, that extremising the action for fixed endpoints of the trajectories gives an equation of motion for the N particles.

- b) With the simplified Lagrangian of the one-dimensional harmonic oscillator

$\mathcal{L}(q, \dot{q}) = \dot{q}^2/2 - q^2/2$, compute and discuss the actions of the following paths $[0, 2\pi) \rightarrow \mathbb{R}$:

- $\gamma_0(t) = 0$
- $\gamma_1(t) = \sin(t)$
- $\gamma_2(t) = \sin(2t)$
- $\gamma_3(t) = t$
- $\gamma_4(t) = e^t$

Exercise 2: Phase Space Part II

We revisit the phase space and study its Hamiltonian structure.

- a) Why is the three-particle elastic collision irresolvable? (“*Physics comes to an end*” does not count as an answer.)
- b) Show that any (positive, scalar) function of the Hamiltonian is a valid choice for a stationary phase-space density.
- c) For a positive energy $E > 0$ in the simplified Lagrangian of the one-dimensional harmonic oscillator $L = \dot{q}^2/2 - q^2/2$, compute the Lebesgue measure of the phase-space set $\{(q, p) \in \mathbb{R}^2 | H(q, p) \leq E\}$, where H is the Hamiltonian function, given as the Legendre transform of L with $\dot{q} \rightarrow p$.

Exercise 3: Symmetry

- a) Show that the gradient transforms like a vector.
- b) Recall (or read up on) *Noether’s Theorem* and convince yourself that for a free classical particle momentum and energy are conserved.

Exercise 4: Probability in Physics

Prove the Stirling formula, following the steps below.

- a) Show that the gamma function generalises the factorial in the following sense:

$$\Gamma(n+1) := \int_0^\infty x^n e^{-x} dx = n!$$

- b) Convince yourself that for large n

$$\int_0^\infty e^{-\frac{(x-n)^2}{2n}} dx \approx \int_{-\infty}^\infty e^{-\frac{(x-n)^2}{2n}} dx.$$

- c) Bring the integrand of the gamma function to the form of a pure exponential function and Taylor expand its exponent around its maximum, dropping all terms higher than second order. Evaluate the so produced integral.