



European Research Council



Existence of the thermodynamic limit for disordered Coulomb quantum systems

Mathieu LEWIN

`mathieu.lewin@math.cnrs.fr`

(CNRS & University of Cergy-Pontoise)

*joint work with X. Blanc (CEA, Paris), based on previous
works with C. Hainzl (Tuebingen) & J.P. Solovej (Copenhagen)*

Spectral Days

Munich, April 11, 2012

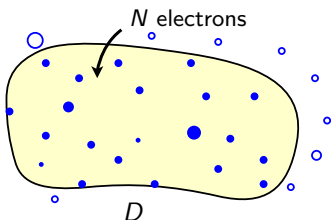
Key references for this talk

	short range	Coulomb
stability of matter $\mathcal{F}(D) \geq C D $	<i>Ruelle & Fisher (66)</i>	<i>Dyson & Lenard (67)</i> <i>Lieb & Thirring (75)</i>
\exists thermo limit, deterministic $\mathcal{F}(D) \sim f D $	<i>Ruelle (63)</i> <i>Fisher (64)</i>	<i>Lebowitz & Lieb (72)</i> <i>Fefferman (85)</i> <i>Hainzl, M.L. & Solovej (09)</i>
\exists thermo limit, stochastic $\mathcal{F}(\omega, D) \sim f D $	<i>Veniaminov (11)</i>	

Main difficulty with Coulomb: quantify screening

Conlon-Lieb-Yau (88), Graf-Schenker (95): lower bounds on classical Coulomb systems when \mathbb{R}^3 split using a tiling

Grand canonical energy



Nuclei: $\mathcal{K} = \{(R, z)\} \subset \mathbb{R}^3 \times (0, \infty)$

$$H_{\mathcal{K}}^N(D) := - \sum_{i=1}^N \Delta_{x_i} + V_{\mathcal{K}, D}(x_1, \dots, x_N)$$

where $-\Delta =$ Dirichlet Laplacian

$$V_{\mathcal{K}, D}(x) = \sum_{i=1}^N \sum_{(R, z) \in \mathcal{K} \cap D} \frac{-z}{|R - x_i|} + \frac{1}{2} \sum_{1 \leq i \neq j \leq N} \frac{1}{|x_i - x_j|} + \frac{1}{2} \sum_{R \neq R' \in \mathcal{K} \cap D} \frac{zz'}{|R - R'|}$$

- ▶ Ground state energy for N electrons:

$$\mathcal{F}_{\mathcal{K}}^N(D) = \inf \left\{ \langle \Psi, H_{\mathcal{K}}^N(D) \Psi \rangle, \Psi \text{ antisymm.}, \|\Psi\|_{L^2(D^N)} = 1 \right\} > -\infty$$

- ▶ Grand canonical energy:

$$\mathcal{F}_{\mathcal{K}}(D) := \inf_{N \geq 0} \mathcal{F}_{\mathcal{K}}^N(D) \quad \left[= \inf_{\bigoplus_{N \geq 0} \bigwedge_1^N L^2(D)} \text{Spec} \bigoplus_{N \geq 0} H_{\mathcal{K}}^N(D) \right]$$

Theorem (Stability of matter)

There exists a constant C , depending on $\max_{(R,z) \in \mathcal{K}} z$ but not on \mathcal{K} , such that

$$\mathcal{F}_{\mathcal{K}}(D) \geq -C|D|$$

for every bounded open set D .

Modern proof based on:

- Lieb-Yau (88) classical electrostatic estimate on many-body Coulomb
 \rightsquigarrow nearest neighbor one-body
- Lieb-Thirring (75) and Sobolev inequalities

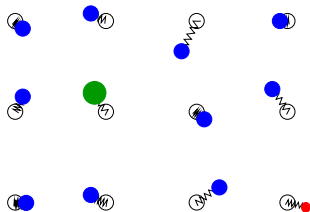
[Sol-06] Solovej, *Stability of Matter*, Encyclopedia of Mathematical Physics, Elsevier (2006).

[LieSei-10] Lieb & Seiringer, *The stability of matter in quantum mechanics*, Cambridge Univ. Press, 2010.

[HaiLewSol-09] Hainzl, M. L. & Solovej, *Advances in Math.*, **221** (2009), pp. 454–487 & 488–546.

Disordered crystal

Typical situation:



$$\mathcal{K}(\omega) = \left\{ (k + \delta_k(\omega), z_k(\omega)), k \in \mathbb{Z}^3 \right\}$$

with δ_k and z_k i.i.d. random variables

Example: $\delta_k \sim$ gaussian and $z_k \sim$ Bernouilli

- \Rightarrow
- ▶ random many-body Schrödinger operator $H^N(\omega, D)$
 - ▶ random grand-canonical energy $\mathcal{F}(\omega, D)$

Ergodicity of nuclei

► General situation:

- translation group $\mathbb{Z}^3 \curvearrowright$ probability space $(\Omega, \mathcal{F}, \mathbb{P})$
- ergodicity: $\tau_k A = A, \forall k \in \mathbb{Z}^3 \Rightarrow \mathbb{P}(A) = 0$ or 1
- $\mathcal{K}(\omega) =$ countable subset of $\mathbb{R}^3 \times [z_{\min}, z_{\max}]$ with

$$\mathcal{K}(\tau_k \omega) = \mathcal{K}(\omega) + k, \quad \text{for all } k \in \mathbb{Z}^3 \text{ and a.s. in } \omega$$

► **Example:** $\Omega = (\Omega_0)^{\mathbb{Z}^3}$, $\delta_k(\omega) = \delta(\omega_k)$, $z_k(\omega) = z(\omega_k)$ and $\tau_k(\omega_\ell) = (\omega_{k+\ell})$

► **Local distribution of nuclei:** with $Q := [0, 1]^3$, define

$$X_0(\omega) := \#(\mathcal{K}(\omega) \cap Q) = \sum_{(R,z) \in \mathcal{K}(\omega) \cap Q} 1$$
$$X_1(\omega) := \sum_{(R,z) \in \mathcal{K}(\omega) \cap Q} \frac{1}{\min_{(R',z') \in \mathcal{K}(\omega) \setminus (R,z)} |R - R'|}$$

Existence of the thermodynamic limit

Theorem (Existence of thermo limit [BlaLew-12])

Assume that

$$\mathbb{E} |X_1|^p < \infty \quad \text{for some } p \geq 2.$$

There exists a cst f such that, for all 'regular' sequences (D_n) with $|D_n| \rightarrow \infty$,

$$\lim_{n \rightarrow \infty} \mathbb{E} \left| \frac{\mathcal{F}(\cdot, D_n)}{|D_n|} - f \right|^q = 0$$

for $q = 1$ if $p = 2$ and all $1 \leq q < p/2$ if $p > 2$.

Rmk.

- Almost sure CV expected as well if (D_n) does not escape to infinity
- Neutrality in average:

$$\lim_{n \rightarrow \infty} \mathbb{E} \left| \frac{N(\cdot, D_n)}{|D_n|} - \mathbb{E} \left(\sum_{(R,z) \in \mathcal{K} \cap Q} z \right) \right|^2 = 0$$

- Same result at finite temperature

Remarks on the assumptions on nuclei

Simple i.i.d. case, ν =law of displacement, charges all =z

- ▶ (Lieb-Yau '88) $\mathbb{E} \frac{\mathcal{F}(\cdot, D)}{|D|} \geq -C + \frac{z^2}{8} \mathbb{E} X_1$, so X_1 must be in $L^1(\Omega)$
- ▶ If ν decays fast, then $\mathbb{E} |X_0|^p < \infty \forall 1 \leq p < \infty$ and $\mathbb{E} |X_1|^p < \infty \forall 1 \leq p < 3$
- ▶ If the support of ν is not compact, then $X_0 \notin L^\infty(\Omega)$
- ▶ If $\nu \geq \varepsilon > 0$ on $B(x, \varepsilon) \cup B(x+k, \varepsilon)$ for some $0 \neq k \in \mathbb{Z}^3$, then $X_1 \notin L^3(\Omega)$

Corollary (Harmonic oscillators)

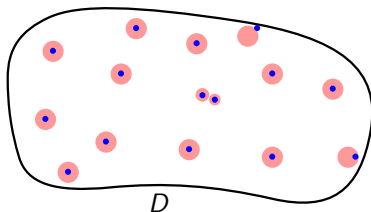
For $\nu(x) = (2\pi\sigma)^{-3/2} \exp(-|x|^2/(2\sigma))$, then

- $\mathbb{E} |X_0|^p < \infty$ for all $1 \leq p < \infty$ and $\mathbb{E} |X_1|^p < \infty$ for all $1 \leq p < 3$;
- X_0 is not bounded and $\mathbb{E} |X_1|^3 = \infty$;
- $\mathbb{E} |\mathcal{F}(\cdot, D)|^3 = \infty$ for all fixed D ;
- the thermo limit exists in $L^q(\Omega)$ for all $1 \leq q < 3/2$.

Strategy of proof

In [HaiLewSol-09], abstract method based on Graf-Schenker inequality

- 1 Stability of matter: **uniform lower bound** on \mathcal{F} ;
- 2 Prove **upper bounds** $\mathbb{E} |\mathcal{F}(\cdot, D)|^q \leq C|D|^q$, for $q \leq p/2$



Screen the nuclei with radial electrons

Radius $\sim \delta_R =$ distance to the nearest nucleus

Cost in kinetic energy: $\sum_{(R,z) \in \mathcal{K}(\omega) \cap D} \delta_R^{-2} \in L^{p/2}(\Omega)$ when $\mathbb{E} |X_1|^p < \infty$

- 3 Prove that periodic $D \mapsto \mathbb{E} \mathcal{F}(\cdot, D)$ satisfies **abstract (deterministic) assumptions of [HaiLewSol-09]**. Requires much more precise upper bounds
 \Rightarrow **thermo limit exists for $\mathbb{E} \mathcal{F}(\cdot, D)$**

- 4 Graf-Schenker inequality together with ergodic thm:

$$\liminf_{n \rightarrow \infty} \frac{\mathcal{F}(\omega, D_n)}{|D_n|} \geq \lim_{n \rightarrow \infty} \frac{\mathbb{E} \mathcal{F}(\cdot, D_n)}{|D_n|} := f, \quad \text{almost surely}$$

- 5 Simple lemma: if (Y_n) bd in $L^1(\Omega)$ is such that $\mathbb{E} Y_n \rightarrow e$, $Y_n \geq C$ and $\liminf Y_n \geq e$, then $\mathbb{E} |Y_n - e| \rightarrow 0$

- 6 The result follows in $L^q(\Omega)$ by interpolation