

**HOMEWORK ASSIGNMENT – WEEK 08-09**

**Hand-in deadline:** Thu 12 June by 12 p.m. in the “MSP” drop box.

**Info:** [www.math.lmu.de/~michel/SS14\\_MSP.html](http://www.math.lmu.de/~michel/SS14_MSP.html)

**Exercise 24.** (Thermodynamic limit of a free Bose gas)

Consider

- a bounded open domain  $\Lambda$  in  $\mathbb{R}^d$ ,  $d \in \mathbb{N}$ ,
- the domains  $\Lambda_L := \{x \in \mathbb{R}^d \mid x/L \in \Lambda\}$ , for every  $L > 0$ ,
- the free particle Hamilton operators  $h_L = -\Delta$  on  $L^2(\Lambda_L)$  with Dirichlet boundary conditions

and use the fact that the spectrum of  $h_1$  consists of a sequence of eigenvalues  $0 \leq \varepsilon_1 < \varepsilon_2 \leq \varepsilon_3 \cdots$  accumulating at infinity according to the Weyl law:

$$\lim_{\varepsilon \rightarrow \infty} \frac{N(\varepsilon)}{\varepsilon^{d/2}} = \text{const},$$

where  $N(\varepsilon)$  is the number of eigenvalues smaller or equal to  $\varepsilon$  (for simplicity we assume  $\varepsilon_1 < \varepsilon_2$ ). In this problem and in the following you are asked to study, for  $\beta > 0$ ,  $0 < z \leq 1$ , the quantity

$$\rho_L(z, \beta) := \frac{1}{|\Lambda_L|} \text{Tr} \frac{z e^{-\beta h_L}}{1 - z e^{-\beta h_L}}$$

when  $L \rightarrow \infty$ . To this aim, fix  $\bar{\rho} > 0$  and define  $z_L$  to be the unique solution to

$$\bar{\rho} = \rho_L(z_L, \beta), \quad z_L := e^{\beta \mu_L}, \quad \mu_L < h_L.$$

In terms of the eigenfunctions  $\psi_N^{(L)}$  of  $h_L$ , namely,  $h_L \psi_N^{(L)} = \varepsilon_N^{(L)} \psi_N^{(L)}$ ,  $N = 1, 2, 3, \dots$ , write

$$\rho_L(z_L, \beta) = \sum_{N=1}^{\infty} \rho_L^{(N)}(z_L, \beta), \quad \rho_L^{(N)}(z_L, \beta) := \frac{1}{|\Lambda_L|} \left\langle \psi_N^{(L)}, \frac{z_L e^{-\beta h_L}}{1 - z_L e^{-\beta h_L}} \psi_N^{(L)} \right\rangle$$

and prove that

$$\lim_{L \rightarrow \infty} \rho_L^{(N)}(z_L, \beta) = 0 \quad \text{when } N > 1.$$

**Exercise 25.** (Follow-up to Exercise 24)

With respect to the assumptions and notation of Exercise 24, prove that

$$\lim_{L \rightarrow \infty} \sum_{n=2}^{\infty} \rho_L^{(n)}(z_L, \beta) = C < \infty$$

where the constant  $C$  is *independent* of  $\bar{\rho}$  (recall that in class it was claimed that the  $C$  amounts exactly to  $\rho_c(\beta)$ ), and thus that

$$\lim_{L \rightarrow \infty} \rho_L^{(1)}(z_L, \beta) > 0.$$

**Exercise 26.** (Bose-Einstein statistics for Weyl's operators)

Recall that in class the KMS states for the free Bose gas were determined in terms of creation and annihilation operators. This required a *regular* representation. In this exercise you are asked to re-do the argument directly for the Weyl's operators, under the assumption that the state is *quasi-free*. To this aim, consider

- the one-particle Hilbert space  $\mathfrak{h}$ ,
- the one-particle Hamiltonian  $h$  on  $\mathfrak{h}$  such that  $h > c\mathbb{1}$  for some  $c > 0$ ,
- the Weyl operator  $W(f)$  as in class, and in particular  $\tau^t(W(f)) = W(e^{ith}f)$ ,  $t \in \mathbb{R}$ ,
- a  $(\tau^t, \beta)$ -KMS state  $\omega$  ( $\beta > 0$ ) which is *quasi-free* and, correspondingly,

$$\xi(t) := \omega(W(-f)\tau^t W(f)).$$

Rewrite  $\xi$  as  $\xi(t) = \omega(W(\eta))F(\eta, t)$  for some  $\eta \in \mathfrak{h}$  that depends on  $t$  and  $f$  and for some  $F(\eta, t)$  that is analytic in  $t$ . Use the analytic continuation of  $F$  to complex-valued arguments to define

$$\xi(i\beta) := \omega(W(\eta))F(\eta, i\beta)$$

and use the KMS condition to determine  $\omega(W(\eta))$ .