## Mathematical Statistical Physics

TMP Programme Munich - spring term 2014

## HOMEWORK ASSIGNMENT - WEEK 11

Hand-in deadline: Thursday 26 June by 12 p.m. in the "MSP" drop box.
Info: www.math.lmu.de/~michel/SS14_MSP.html

Exercise 30. (The lattice BCS model.)
Consider, among the class of quantum spin systems studied in Exercise 29 in the mean field regime, the special case of a one-dimensional spin- $\frac{1}{2}$ model where the Hamiltonian $H_{\Lambda}$ is given, with respect to the notation of Exercise 29, by

$$
d=1, \quad s=\frac{1}{2}, \quad A=\left(\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right), \quad B=-\lambda\left(C^{*} \otimes C+C \otimes C^{*}\right), \quad C=\left(\begin{array}{ll}
0 & 1 \\
0 & 0
\end{array}\right), \quad \lambda \geqslant 0 .
$$

(i) Prove that a generic $2 \times 2$ density matrix $\rho$ can be parametrised as follows:

$$
\rho=\left(\begin{array}{cc}
\frac{1}{2}+r & \mu \\
\bar{\mu} & \frac{1}{2}-r
\end{array}\right), \quad r \in \mathbb{R}, \quad \mu \in \mathbb{C}, \quad r^{2}+|\mu|^{2} \leqslant \frac{1}{4} .
$$

(ii) Compute the reduced Hamiltonian $H_{\rho}$ of the model relative to a density matrix $\rho$ parametrised as in (i).
(iii) Compute $e^{-\beta H_{\rho}}, \beta>0$.
(iv) Write and solve the gap equation (see Exercise 29(ii)) for this model. Plot the parameter $r$ and $|\mu|$ of the solutions as a function of $\beta$.
(v) Which of the solutions found in (iv) actually maximises $f_{\beta}$ ?

Exercise 31. (Application of RG: a quantum flute - or a bosonic string)
Consider a cigar-shaped tube of length $\pi$ and assume that the pressure inside it can be modelled by a function $p(t, x)$ of the time $t$ and of the coordinate $x$ along the tube. Consider the classical Hamiltonian

$$
H=\int_{0}^{\pi} \mathrm{d} x\left(\left(\frac{\partial p}{\partial t}\right)^{2}+\left(\frac{\partial p}{\partial x}\right)^{2}\right) .
$$

(i) Assume Dirichlet boundary conditions $p(t, 0)=p(t, \pi)=0$. Re-write $H$ in terms of the Fourier transform of $p(t, x)$ with respect to $x$ and prove that $H$ takes this way the form of an infinite sum of classical harmonic oscillators of mass $\frac{1}{2}$,

$$
H=\sum_{k=1}^{\infty} h_{k}, \quad h_{k}:=\left(\frac{\mathrm{d} p_{k}}{\mathrm{~d} t}\right)^{2}+\omega_{k} p_{k}^{2},
$$

of which you have to specify the frequency $\omega_{k}$.
(ii) Consider the quantisation of the Hamiltonian $H$ obtained in (i), consisting of the replacement of each mode $p_{k}$ with the operator of multiplication times $x_{k}$ and correspondingly of $\frac{\mathrm{d} p_{k}}{\mathrm{~d} t}$ with the differential operator $-\mathrm{i} \frac{\partial}{\partial x_{k}}$. Let $E_{k}$ be the ground state energy of the harmonic oscillator $h_{k}$. Prove that $\sum_{k=1}^{\infty} E_{k}=+\infty$.
(iii) To cure the divergence found in (ii), assume that $H$ can be regularised to some $H_{\text {reg }}$ in such a way that each harmonic oscillator $h_{k}$ has a ground state damped by a factor $e^{-a / \lambda_{k}}$, where $a$ is a reference distance (say, the typical inter-atomic distance) and $\lambda_{k}$ is the wave length of the mode $k$. Correspondingly, consider $\widetilde{E}:=\sum_{k=1}^{\infty} E_{k} e^{-a / \lambda_{k}}$. Prove the following asymptotics

$$
\widetilde{E}=\frac{1}{a^{2}}-\frac{1}{12}+O(a) \quad \text { as } a \rightarrow 0 .
$$

(iv) An immediate consequence of (iii) is that re-normalising $H_{\text {reg }}$ by

$$
H_{\mathrm{reg}} \mapsto H^{\prime}:=H_{\mathrm{reg}}-\int_{0}^{\pi} \mathrm{d} x \frac{1}{a^{2}},
$$

namely adding a contribution which is $p$-independent (thus, not affecting the equation of motion) and preserves locality, the ground state energy of $H^{\prime}$ stays finite as $a \rightarrow 0$. Discuss the value of this ground state in relation to the quantity $\zeta(-1)$, where $\zeta$ is the Riemann zeta function.

Exercise 32. For each of the 10 questions below answer YES or NO and provide a brief explanation.
32.1 Is there a unique KMS state for an infinite block of iron at room temperature?

## $\square$ YES $\square$ NO

32.2 Consider the $C^{*}$-algebra $\mathcal{A}=\mathcal{L}\left(L^{2}(\mathbb{R})\right)$ and the time evolution $\tau_{t}(A)=e^{\mathrm{it} H} A e^{-\mathrm{i} t H}, t \in \mathbb{R}$, where $H$ is the Hamiltonian of the 1D harmonic oscillator, namely $(H f)(x):=-f^{\prime \prime}(x)+x^{2} f(x)$ on the domain $\mathcal{D}=\left\{f \in L^{2}(\mathbb{R}) \mid-f^{\prime \prime}+x^{2} f \in L^{2}(\mathbb{R})\right\}$. Is $\tau_{t}$ asymptotically abelian?

- YES
32.3 With the usual meaning of the symbols from class and homework, consider the 1D Ising model at non-zero temperature $T$ and with magnetic field $B$. Is it true that at the Renormalisation Group limit point one has $\left\langle S_{1} S_{2}\right\rangle=\left\langle S_{1}\right\rangle\left\langle S_{2}\right\rangle$ ?
- YES
32.4 With the usual meaning of the symbols from class and homework, consider the 3D Ising model. If you plot the logarithm of the spontaneous magnetization as a function of $\log \frac{\left(T-T_{c}\right)}{T_{c}}$ for $T$ approaching $T_{c}$ from below, do you get a straight line?
$\square$ YES $\square$ NO
32.5 Consider the spin- $\frac{1}{2} 2 \mathrm{D}$ Heisenberg model on $\mathbb{Z}^{2}$ and, for every $L \in \mathbb{Z}$, consider the finite sub-lattice $\Lambda_{L}$ consisting of the square of size $L$ centred at the origin. Is it possible to find for every $L$ a finite sub-lattice $\Lambda_{L}^{\prime} \supsetneq \Lambda_{L}$ and a configuration on the whole $\mathbb{Z}^{2}$ with all spins up inside $\Lambda_{L}$ and all spins down outside $\Lambda_{L}^{\prime}$ such that $\lim _{L \rightarrow \infty} \mathcal{E}(L)$ is finite? $(\mathcal{E}(L):=$ energy of the configuration relative to the couple $\Lambda_{L}, \Lambda_{L}^{\prime}$.)
$\square$ YES
- NO
32.6 Consider an "Ideal Bose Gas consisting of one-body harmonic oscillators", i.e., with respect to Fock-space discussion of the Ideal Bose Gas done in class and homework, assume that the one-body Hamiltonian is a harmonic oscillator of frequency $\omega$. Does BEC occur in the limit $\omega \rightarrow 0$ ?
- YES $\square \mathrm{NO}$
32.7 Consider a process in which at each step $t$ a particle jumps with probability $\frac{1}{3}$ from a corner of a given cube to one of the three neighbouring corners. Is this a Markov process that converges to a probability distribution in the limit $t \rightarrow \infty$ ?


## $\square$ YES $\square$ NO

32.8 Let $p, q \in(0,1)$ such that $p+q<1$. Consider four sites $S_{1}, S_{2}, S_{3}$, and $S_{4}$, and a process in which at each step $t$ a particle jumps from one site to another, say $S_{j} \rightarrow S_{k}$, with this prescription:

- $S_{j} \rightarrow S_{j}$ with probability $p(j=1,2,3,4)$,
- $S_{1} \rightarrow S_{2}$ or $S_{4} \rightarrow S_{3}$ with probability $1-p$,
- $S_{2} \rightarrow S_{1}$ or $S_{3} \rightarrow S_{4}$ with probability $1-p-q$,
- $S_{2} \rightarrow S_{3}$ or $S_{3} \rightarrow S_{2}$ with probability $q$.

According to what was discussed in class, this gives in fact a Markov process that converges to a limiting distribution $\rho_{\infty}$ (because of aperiodicity and irreducibility) and, if $\rho(t)$ is the distribution at time $t$, the displacement $\rho(t)-\rho_{\infty}$ vanishes as $e^{-\mu t}$ for some $\mu>0$. Is it true that $\mu=O(q)$ as $q \rightarrow 0$ ?
$\square \mathrm{YES} \square \mathrm{NO}$
32.9 Consider a $C^{*}$-dynamical system $\left(\mathcal{A}, \tau^{t}\right)$, a $\left(\tau^{t}, \beta\right)$-KMS state $\omega(\beta>0)$ on $\mathcal{A}$, and a *-automorphism $\alpha$ on $\mathcal{A}$. Is it true that, if $\omega$ is $\alpha$-invariant, so is $\tau^{t}$ ?
$\square$ YES $\square$ NO
32.10 Consider the $C^{*}$-algebra $\mathcal{A}=\mathcal{M}(3 \times 3, \mathbb{C})$, the state $\rho=\left(\begin{array}{ccc}\frac{1}{2} & 0 & 0 \\ 0 & \frac{1}{4} & 0 \\ 0 & 0 & \frac{1}{4}\end{array}\right)$, a dynamics $\tau^{t}$ on $\mathcal{A}$ such that $\rho$ is a $\left(\tau^{t}, 1\right)$-KMS state, and the manifold

$$
\mathcal{M}_{\tau}:=\left\{\widetilde{\rho} \mid \widetilde{\rho} \text { is a state on } \mathcal{A} \text { and } \widetilde{\rho} \circ \tau^{t}=\widetilde{\rho} \forall t \in \mathbb{R}\right\}
$$

of all $\tau^{t}$-invariant states on $\mathcal{A}$. Is it true that $\operatorname{dim}_{\mathbb{R}} \mathcal{M}_{\tau}=4$ ?
$\square \mathrm{YES} \square \mathrm{NO}$

